

NASA TM X-63011

LONG-PERIOD VARIATIONS IN THE ECCENTRICITY AND ARGUMENT OF PERIGEE OF NIMBUS 2

THEODORE L. FELSENTREGER

MAY 1967



GODDARD SPACE FLIGHT CENTER
GREENBELT, MARYLAND

EF No. 602(C)	<u>N 68 - 10257</u>	<u> </u>
	(ACCESSION NUMBER)	(THRU)
	<u>17</u>	<u>1</u>
	(PAGES)	(CODE)
	<u>TMX-63011</u>	<u>30</u>
	(NASA CR OR TMX OR AD NUMBER)	(CATEGORY)

X-547-67-234

LONG-PERIOD VARIATIONS IN THE ECCENTRICITY
AND ARGUMENT OF PERIGEE OF NIMBUS 2

Theodore L. Felsentreger

May 1967

Goddard Space Flight Center
Greenbelt, Maryland

LONG-PERIOD VARIATIONS IN THE ECCENTRICITY AND ARGUMENT OF PERIGEE OF NIMBUS 2

Theodore L. Felsentreger

Summary

Long-period variations have been observed in the mean values of eccentricity and argument of perigee (published by the Goddard Space Flight Center) for Nimbus 2. It is shown that the use of a more recent determination of the earth zonal harmonic coefficients, plus additional first-order earth oblateness perturbations (for small eccentricity satellites) not included in the mathematical model for orbit determination, explains these variations.

LONG-PERIOD VARIATIONS IN THE ECCENTRICITY AND ARGUMENT OF PERIGEE OF NIMBUS 2

INTRODUCTION

In two earlier papers (References 1 and 2), the problem of small divisors in the case of low eccentricity satellites was analyzed at some length. Eccentricity was found to occur as a divisor in analytic expressions (both long and short period) for the perturbations, caused by the non-sphericity of the earth, in eccentricity, mean anomaly, and argument of perigee. Formulas for these perturbations were presented, and contain terms involving up to J_2^3 and $(J_1/J_2)^3$, where J_1 is any odd zonal harmonic. In addition, the long-period terms were found to explain satisfactorily the observed variations in eccentricity and argument of perigee for the satellites Alouette 1 and Tiros 8.

Similar long-period variations in eccentricity and argument of perigee were observed in the case of Nimbus 2, another low eccentricity satellite. It will be shown here that the long-period expressions in References 1 and 2 are adequate representations of these perturbations.

LUNAR AND SOLAR EFFECTS

Mean values for the eccentricity and argument of perigee, developed at the Goddard Space Flight Center, were first corrected for lunar and solar gravitational effects and the effects caused by solar radiation pressure and lunar and solar tides; the formulas used to compute these perturbations appear in Reference 3 and 4. For the most part, these effects were fairly small, except for three near-resonant solar terms in the argument of perigee perturbations. However, since these terms all had periods in excess of 10^4 days, their effects over the time interval studied would appear to be secular and therefore wouldn't affect the analysis of the periodic variations.

The corrected values of eccentricity and argument of perigee appear in Tables 1 and 2 as e_c'' and g_c'' , respectively.

LONG-PERIOD ZONAL HARMONIC PERTURBATIONS

Analytic formulas for the major long-period zonal harmonic effects on the eccentricity and argument of perigee for a low eccentricity satellite appear in References 1 and 2; they are reproduced here for convenience:

$$\begin{aligned}
 \Delta e = & - \left[Q - \frac{Q^3}{8e_1^2} - \frac{J_2^2}{64e_1^2 a''^4} (28 - 101 \cos^2 i'' + 218 \cos^4 i'') Q \right] \cos \bar{\theta} \\
 & - \left[\frac{Q^2}{4e_1} + \frac{9J_2^2}{64e_1 a''^4} \sin^2 i'' (1 - 3 \cos^2 i'') \right] \cos 2 \bar{\theta} \\
 & - \left[\frac{Q^3}{8e_1^2} + \frac{27J_2^2}{128e_1^2 a''^4} \sin^2 i'' (1 - 3 \cos^2 i'') Q \right] \cos 3 \bar{\theta} \\
 \Delta g = & \left(\frac{Q}{e_1} + \frac{Q^3}{4e_1^3} \right) \sin \bar{\theta} + \left[\frac{Q^2}{2e_1^2} + \frac{9J_2^2 \sin^2 i'' (1 - 3 \cos^2 i'')}{32e_1^2 a''^4} \right] \sin 2 \bar{\theta} \\
 & + \left[\frac{Q^3}{3e_1^3} + \frac{9J_2^2 \sin^2 i'' (1 - 3 \cos^2 i'')}{16e_1^3 a''^4} Q \right] \sin 3 \bar{\theta}, \quad (1)
 \end{aligned}$$

where a'' , e_1 , and i'' are semi-major axis, eccentricity, and inclination constants, respectively. In addition,

$$\bar{\theta} = \text{constant} + \text{secular term} \simeq g'' + \frac{\pi}{2}$$

g'' = argument of perigee

$$Q = \frac{M}{N},$$

where

$$\begin{aligned}
N = & - \frac{3J_2}{4a''^3\sqrt{a''}} (1 - 5 \cos^2 i'') + \frac{3J_2^2}{64a''^5\sqrt{a''}} (7 - 114 \cos^2 i'' + 395 \cos^4 i'') \\
& - \frac{15J_4}{32a''^5\sqrt{a''}} (3 - 36 \cos^2 i'' + 49 \cos^4 i'') + \dots \\
M = & \frac{3J_3 \sin i''}{8a''^4\sqrt{a''}} (1 - 5 \cos^2 i'') + \frac{15J_5 \sin i''}{32a''^6\sqrt{a''}} (1 - 14 \cos^2 i'' + 21 \cos^4 i'') \\
& + \frac{105J_7 \sin i''}{1024a''^8\sqrt{a''}} (5 - 135 \cos^2 i'' + 495 \cos^4 i'' - 429 \cos^6 i'') \\
& + \frac{315J_9 \sin i''}{4096a''^{10}\sqrt{a''}} (7 - 308 \cos^2 i'' + 2002 \cos^4 i'' - 4004 \cos^6 i'' + 2431 \cos^8 i'') \\
& + \frac{3465J_{11} \sin i''}{131072a''^{12}\sqrt{a''}} (21 - 1365 \cos^2 i'' + 13650 \cos^4 i'' - 46410 \cos^6 i'' \\
& + 62985 \cos^8 i'' - 29393 \cos^{10} i'') + \dots
\end{aligned} \tag{2}$$

However, the terms $-Q \cos \bar{\theta}$ (for Δe) and $(Q/e_1) \sin \bar{\theta}$ (for Δg) have already been used in the mathematical model for orbit determination, at least for the zonal harmonics J_2 , J_3 , J_4 , and J_5 . The following values were used for J_3 and J_5 :

$$J_3 = -2.285 \times 10^{-6}, \quad J_5 = -0.232 \times 10^{-6}.$$

It is suggested that Kozai's 1964 determination of the harmonic coefficients (see Reference 5) form a better set — they are

$$J_2 = 1.082645 \times 10^{-3}$$

$$J_3 = -2.546 \times 10^{-6}$$

$$J_4 = -1.649 \times 10^{-6}$$

$$J_5 = -0.210 \times 10^{-6}$$

$$J_7 = -0.333 \times 10^{-6}$$

$$J_9 = -0.053 \times 10^{-6}$$

$$J_{11} = 0.302 \times 10^{-6}.$$

It is further suggested that the observed variations in e and g would reflect the differences between the more accurate J_3 , J_5 values and the values which were actually used in the $-Q \cos \bar{\theta}$ and $(Q/e_1) \sin \theta$ terms.

Therefore, setting

$$\Delta J_3 = -2.546 \times 10^{-6} - (-2.285 \times 10^{-6}) = -0.261 \times 10^{-6}$$

$$\Delta J_5 = -0.210 \times 10^{-6} - (-0.232 \times 10^{-6}) = 0.22 \times 10^{-7}$$

$$\begin{aligned} \Delta M = & \frac{3\Delta J_3 \sin i''}{8a''^4 \sqrt{a''}} (1 - 5 \cos^2 i'') + \frac{15\Delta J_5 \sin i''}{32a''^6 \sqrt{a''}} (1 - 14 \cos^2 i'' + 21 \cos^4 i'') \\ & + \frac{105J_7 \sin i''}{1024a''^8 \sqrt{a''}} (5 - 135 \cos^2 i'' + 495 \cos^4 i'' - 429 \cos^6 i'') \\ & + \dots \end{aligned}$$

$$\Delta Q = \frac{\Delta M}{N}, \quad (3)$$

the expressions to be compared with the observed variations are

$$\begin{aligned} \Delta e = & - \left[\Delta Q - \frac{Q^3}{8e_1^2} - \frac{J_2^2}{64e_1^2 a''^4} (28 - 101 \cos^2 i'' + 218 \cos^4 i'') Q \right] \cos \bar{\theta} \\ & - \left[\frac{Q^2}{4e_1} + \frac{9J_2^2}{64e_1 a''^4} \sin^2 i'' (1 - 3 \cos^2 i'') \right] \cos 2\bar{\theta} \\ & - \left[\frac{Q^3}{8e_1^2} + \frac{27J_2^2}{128e_1^2 a''^4} \sin^2 i'' (1 - 3 \cos^2 i'') Q \right] \cos 3\bar{\theta} \end{aligned}$$

$$\Delta g = \left(\frac{\Delta Q}{e_1} + \frac{Q^3}{4e_1^3} \right) \sin \bar{\theta} + \left[\frac{Q^2}{2e_1^2} + \frac{9J_2^2 \sin^2 i'' (1 - 3 \cos^2 i'')}{32e_1^2 a''^4} \right] \sin 2\bar{\theta} \\ + \left[\frac{Q^3}{3e_1^3} + \frac{9J_2^2 \sin^2 i'' (1 - 3 \cos^2 i'')}{16e_1^3 a''^4} Q \right] \sin 3\bar{\theta} . \quad (4)$$

Using

$$a'' = 1.1782076 \text{ earth radii}$$

$$e_1 = .00560$$

$$i'' = 100^\circ 306$$

and Kozai's 1964 determination of the zonal harmonic coefficients, Equations (4) become

$$\Delta e = -(1.0923971 \times 10^{-4}) \cos \bar{\theta} - (0.63515970 \times 10^{-4}) \cos 2\bar{\theta} + (0.85403736 \times 10^{-5}) \cos 3\bar{\theta}$$

$$\Delta g = 1^\circ 3454747 \sin \bar{\theta} + 1^\circ 2997132 \sin 2\bar{\theta} + 0^\circ 23301303 \sin 3\bar{\theta}. \quad (5)$$

OBSERVED VARIATIONS

The values of e_c'' and g_c'' were then fit, by means of least squares, to the following models:

$$e_c'' = e_0 + \sum_{n=1}^9 A_n \cos n \bar{\theta} \\ g_c'' = g_0 + \dot{g}_c'' (t - t_0) + \sum_{n=1}^9 B_n \sin n \bar{\theta}, \quad (6)$$

where e_0 , g_0 , A_n , and B_n are constants, \dot{g}_c'' is the secular motion of g_c'' , $t - t_0$ is the elapsed time in days, and

$$\bar{\theta} = 51^\circ 34033 - (2^\circ 3536297/\text{day}) (t - t_0).$$

The constant $51^{\circ}34'03.3$ is merely the first value of g_c'' plus 90° , and the secular motion of $\bar{\theta}$ was computed from the expression for dg''/dt appearing in Reference 6. The results are

$$\begin{aligned} e_c'' = & 0.55936191 \times 10^{-2} - (1.200321 \times 10^{-4}) \cos \bar{\theta} - (0.648964 \times 10^{-4}) \cos 2\bar{\theta} \\ & + (0.44700 \times 10^{-5}) \cos 3\bar{\theta} + (1.5399 \times 10^{-6}) \cos 4\bar{\theta} + (1.7144 \times 10^{-6}) \cos 5\bar{\theta} \\ & - (0.9668 \times 10^{-6}) \cos 6\bar{\theta} + (0.34168 \times 10^{-5}) \cos 7\bar{\theta} + (1.9838 \times 10^{-6}) \cos 8\bar{\theta} \\ & + (0.8763 \times 10^{-6}) \cos 9\bar{\theta} \end{aligned}$$

$$\begin{aligned} g_c'' = & 679^{\circ}62'28.5 - (2^{\circ}35'42.110/\text{day}) (t - t_0) + 1^{\circ}30'07.803 \sin \bar{\theta} \\ & + 0^{\circ}8'19.2370 \sin 2\bar{\theta} - 0^{\circ}0'20.2026 \sin 3\bar{\theta} + 0^{\circ}0'19.6037 \sin 4\bar{\theta} \\ & + 0^{\circ}0'09.9007 \sin 5\bar{\theta} + 0^{\circ}0'44.1686 \sin 6\bar{\theta} + 0^{\circ}0'03.4652 \sin 7\bar{\theta} \\ & + 0^{\circ}0'37.9263 \sin 8\bar{\theta} + 0^{\circ}0'27.9919 \sin 9\bar{\theta}. \end{aligned} \tag{7}$$

A comparison between Equations (5) and (7) indicates that the major observable periodic variations in e and g (the $\bar{\theta}$ and $2\bar{\theta}$ terms) are explainable by the theory. There is somewhat of a discrepancy between the amplitudes of the $\sin 2\bar{\theta}$ terms. However, it should be noted (see Equations (4)) that the theoretical amplitude is sensitive to the value of e_1 used; use of a larger value would bring closer agreement.

Figures 1 and 2 show the closeness of the least squares fits.

ACKNOWLEDGMENT

The author wishes to thank Mr. Wilbur B. Huston of the Nimbus Project Office for drawing attention to this problem, and also for several helpful discussions.

REFERENCES

1. Felsentreger, Theodore L., and Victor, Eric L., "On the Long Period Perturbations in the Motion of Small Eccentricity Satellites," GSFC X-547-66-577, December, 1966.
2. Felsentreger, Theodore L., and Steinberg, Ellen L., "On the Perturbations of Small Eccentricity Satellites," GSFC X-547-67-102, March, 1967.
3. Murphy, James P., and Felsentreger, Theodore L., "Analysis of Lunar and Solar Effects on the Motion of Close Earth Satellites," NASA TN D-3559, August, 1966.
4. Fisher, David, and Felsentreger, Theodore L., "Effects of the Solar and Lunar Tides on the Motion of an Artificial Satellite," GSFC X-547-66-560, November, 1966.
5. Kozai, Y., "New Determination of Zonal Harmonics Coefficients of the Earth's Gravitational Potential," SIAO Special Report No. 165, November 2, 1964.
6. Brouwer, D., "Solution of the Problem of Artificial Satellite Theory Without Drag," Astronomical Journal 64 (9), pp. 378-397, November, 1959.

Table 1. Eccentricity of Nimbus 2

$t - t_0(\text{days})$	$e_c'' \times 10^2$	$\left(e_0 + \sum_{n=1}^9 A_n \cos n\bar{\theta} \right) \times 10^2$	
0	.55332889	.55309715	
7	.54728658	.54689101	
14	.54168522	.54263530	
21	.54233060	.54216936	
28	.54606277	.54230895	
35	.54525536	.54553649	
42	.54996346	.55172237	
49	.55606657	.55742556	
55	.55957010	.56257593	
62	.56235129	.56728237	
69	.56480941	.56868472	
76	.56463475	.56837489	
83	.56371704	.56742564	
90	.56290166	.56565369	
97	.56267565	.56413156	
104	.56470174	.56492539	
111	.56764970	.56686812	
118	.57106100	.56810100	
125	.57244757	.56866868	
132	.57036198	.56814157	
139	.56764492	.56465560	
146	.56238443	.55874767	
153	.55574878	.55306014	
160	.54816344	.54685197	
167	.54255936	.54262321	
174	.54007409	.54216969	
181	.54265645	.54231507	
188	.54266910	.54557165	
195	.55178840	.55176106	
202	.55480257	.55746235	
209	.56466682	.56344980	
216	.56788385	.56768853	
223	.56750003	.56869821	
230	.56620046	.56827280	
237	.56768515	.56721607	
244	.56621362	.56534977	
251	.56548537	.56408524	
258	.56634020	.56520973	
264	.56908258	.56687891	
271	.57008731	.56810638	
278	.57087450	.56867023	
285	.56981932	.56813067	
292	.56534614	.56462217	
299	.55772803	.55870967	
306	.55368415	.55302309	
$e_0 = .0055936191$	$A_1 = -.0001200321$	$A_2 = -.0000648964$	$A_3 = .0000044700$
$A_4 = .0000015399$	$A_5 = .0000017144$	$A_6 = -.0000009668$	$A_7 = .0000034168$
$A_8 = .0000019838$	$A_9 = .0000008763$	$\bar{\theta} = 51:34033 - (2:3536297/\text{day}) (t - t_0)$	
$t - t_0 = \text{days since } 5/22/66, 0 \text{ hrs. U.T.}$			

Table 2. Argument of Perigee of Nimbus 2

$t - t_0$ (days)	g_c'' (deg.)	$g_c'' - [g_0 + \dot{g}_c''(t - t_0)]$ (deg.)	$\sum_{n=1}^9 B_n \sin n \bar{\theta}$ (deg.)
0	681.34033	1.71748	1.81008
7	664.56280	1.41943	1.42397
14	647.43303	0.76913	0.98474
21	630.18123	-0.00319	0.12833
28	612.91141	-0.79353	-0.84045
35	596.17881	-1.04665	-1.32491
42	579.25093	-1.49506	-1.74587
49	562.60975	-1.65676	-1.81638
55	548.70462	-1.43662	-1.60563
62	532.58137	-1.08040	-1.22130
69	516.59107	-0.59122	-0.63771
76	500.47859	-0.22422	-0.25701
83	484.09840	-0.12494	0.00321
90	467.60543	-0.13843	0.15816
97	450.82827	-0.43611	0.03784
104	434.13663	-0.64828	-0.14006
111	417.82345	-0.48198	-0.08924
118	401.75990	-0.06605	0.16244
125	385.67054	0.32406	0.46239
132	369.87252	1.00552	1.01109
139	354.07590	1.68838	1.49166
146	337.83500	1.92696	1.77110
153	321.44596	2.01739	1.80866
160	305.04411	2.09502	1.42115
167	287.97163	1.50202	0.98117
174	270.32661	0.33647	0.12129
181	252.80550	-0.70516	-0.84488
188	236.12142	-0.90976	-1.32760
195	218.76501	-1.78669	-1.74801
202	202.22859	-1.84364	-1.81532
209	186.04263	-1.55012	-1.56182
216	169.99133	-1.12194	-1.13988
223	153.95519	-0.67861	-0.56128
230	137.79547	-0.35885	-0.21829
237	121.45899	-0.21585	0.04051
244	105.04672	-0.14865	0.15577
251	88.51341	-0.20248	0.00728
258	71.94073	-0.29568	-0.15223
264	57.97534	-0.13581	-0.08798
271	41.83007	0.19840	0.16406
278	25.69696	0.54477	0.46502
285	9.76031	1.08759	1.01492
292	- 6.29829	1.50843	1.49380
299	-22.50038	1.78586	1.77253
306	-38.95179	1.81392	1.80722
$g_0 = 679^{\circ}62285 \quad g_c'' = -2^{\circ}3542110/\text{day} \quad B_1 = 1^{\circ}3007803 \quad B_2 = 0^{\circ}8192370$			
$B_3 = -0^{\circ}0202026 \quad B_4 = 0^{\circ}0196037 \quad B_5 = 0^{\circ}0099007 \quad B_6 = 0^{\circ}0441686$			
$B_7 = 0^{\circ}0034652 \quad B_8 = 0^{\circ}0379263 \quad B_9 = 0^{\circ}0279919$			
$\bar{\theta} = 51^{\circ}34033 - (2^{\circ}3536297/\text{day})(t - t_0)$			

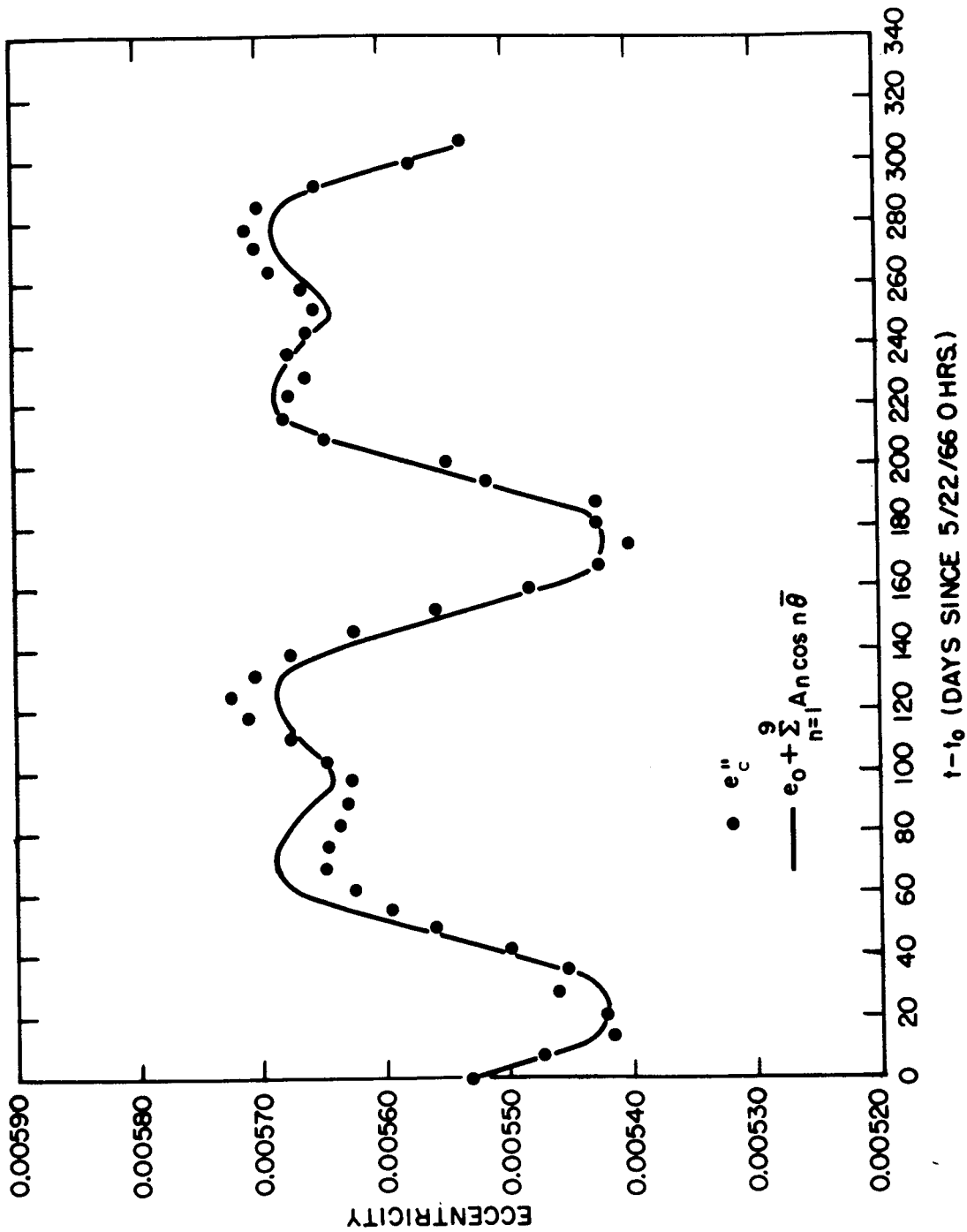


Figure 1—Eccentricity of Nimbus 2.

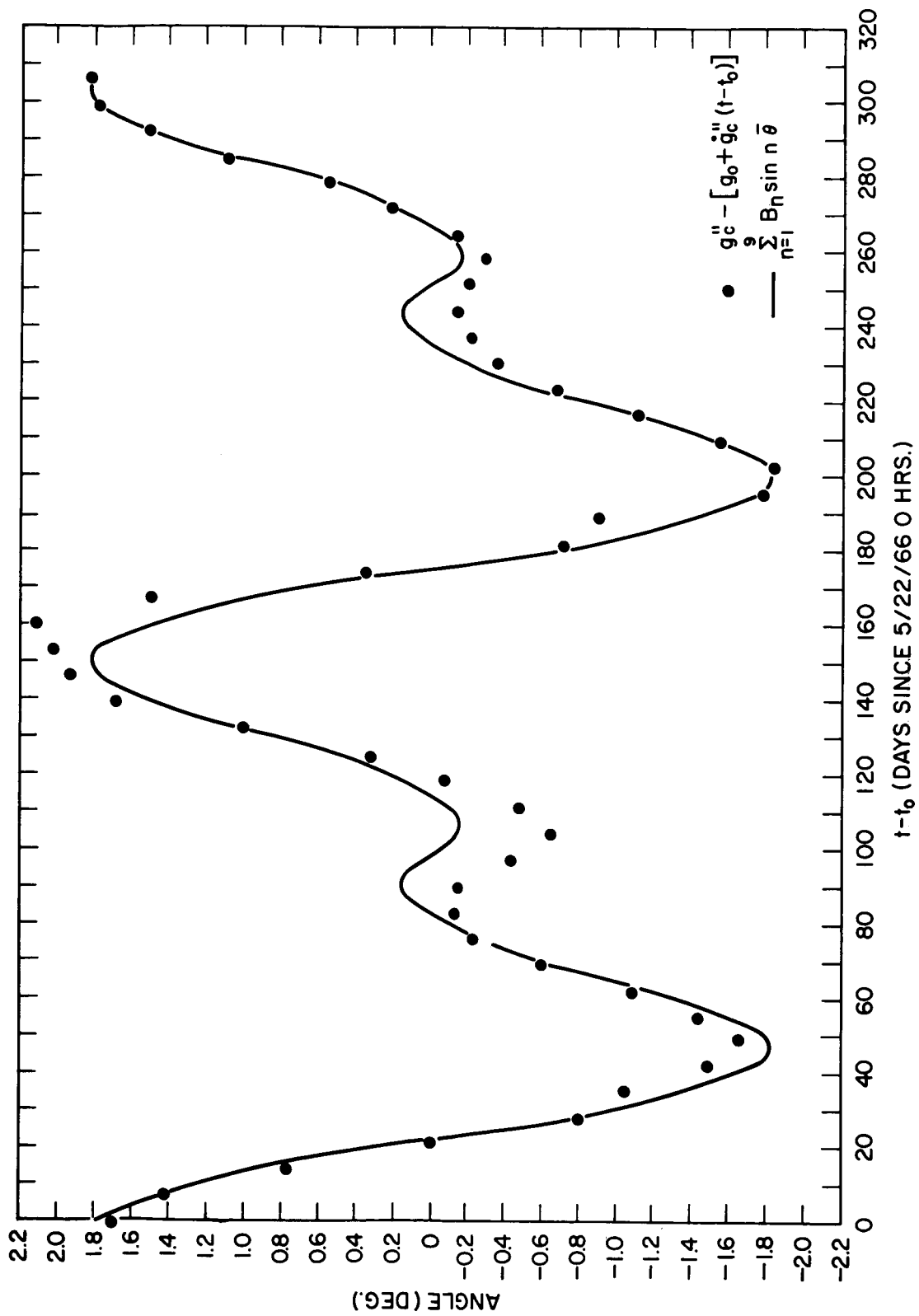


Figure 2—Argument of Perigee of Nimbus 2.